# A Reduced-Order Model for the Dynamics of a Flexible Solar Sail 

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- designing a feedback control to stabilize the solarcraft about the desired maneuver


## INTRODUCTION

Flexible solarcraft motion involves 3 fundamental disciplines:

INTRODUCTION

## Dynamics

INTRODUCTION


INTRODUCTION


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We consider a square solar sail.

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We assume the sail film is connected to the booms continuously.

## EQUATIONS OF MOTION



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Body axes $x y z$ are obtained from the inertial axes $X Y Z$ :

1. Rotation $\theta$ about $Y$ to the axes $x_{1} y_{1} z_{1}$
2. Rotation $\phi$ about $x_{1}$ to the axes $x_{2} y_{2} z_{2}$
3. Rotation $\psi$ about $z_{2}$ to the body axes $x y z$
$\phi, \theta$ and $\psi$ are the Euler angles.

## EQUATIONS OF MOTION

The velocity of the origin $o$ of the body axes $\mathbf{v}=\left[\begin{array}{lll}v_{x} & v_{y} & v_{z}\end{array}\right]^{T}$.

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\mathbf{v}=C(\phi, \theta, \psi) \dot{\mathbf{R}}
$$

In body axes components

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$$
\mathbf{v}=C(\phi, \theta, \psi) \dot{\mathbf{R}} \quad \rightarrow \quad \mathbf{R}=C^{T}(\phi, \theta, \psi) \mathbf{v} \text { inertial axes components }
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\begin{aligned}
& \mathbf{v}=C(\phi, \theta, \psi) \dot{\mathbf{R}} \rightarrow \quad \dot{\mathbf{R}}=C^{T}(\phi, \theta, \psi) \mathbf{v} \\
& \mathbf{\omega}=E(\phi, \psi) \dot{\boldsymbol{\theta}}
\end{aligned}
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\mathbf{v}=C(\phi, \theta, \psi) \dot{\mathbf{R}} & \rightarrow & \dot{\mathbf{R}}=C^{T}(\phi, \theta, \psi) \mathbf{v} \\
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$\mathbf{v}$ and $\boldsymbol{\omega}$ are referred to as quasi-velocities.

## EQUATIONS OF MOTION

Local axes for Control Vane 1-4 are obtained from xyz:

1. Rotation $\alpha_{i}$ about $z$ to the axes $x_{i}^{\prime} y_{i}^{\prime} z_{i}^{\prime}$
2. Rotation $\beta_{i}$ about $x_{i}^{\prime}$ to the axes $x_{i}^{\prime \prime} y_{i}^{\prime \prime} z_{i}^{\prime \prime}$
3. Rotation $\delta_{i}$ about $y_{i}^{\prime \prime}$ to the local axes $\eta_{i} \gamma_{i} \kappa_{i}$
for $i=1,2,3,4$.
For the square solar sail, $\alpha_{1}=225^{\circ}, \alpha_{2}=45^{\circ}, \alpha_{3}=315^{\circ}, \alpha_{4}=135^{\circ}$.

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\frac{d}{d t}\left(\frac{\partial \mathcal{L}}{\partial \boldsymbol{\omega}}\right)+\tilde{v} \frac{\partial \mathcal{L}}{\partial \mathbf{v}}+\tilde{\omega} \frac{\partial \mathcal{L}}{\partial \boldsymbol{\omega}}-\left(E^{T}\right)^{-1} \frac{\partial \mathcal{L}}{\partial \boldsymbol{\theta}} & =\mathbf{M} \\
\frac{\partial}{\partial t}\left(\frac{\partial \hat{\mathcal{L}}}{\partial \dot{\mathbf{u}}}\right)-\frac{\partial \hat{\mathcal{L}}}{\partial \mathbf{u}}+\frac{\partial \hat{\mathcal{R}}}{\partial \dot{\mathbf{u}}} & =\hat{\mathbf{U}}
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$\begin{aligned} \begin{array}{l}\text { Force Equations: } \\ 3 \text { nonlinear ODEs }\end{array} & \begin{array}{r}\frac{d}{d t}\left(\frac{\partial \mathcal{L}}{\partial \mathbf{v}}\right)+\tilde{\omega} \frac{\partial \mathcal{L}}{\partial \mathbf{v}}-C \frac{\partial \mathcal{L}}{\partial \mathbf{R}} \\ \frac{d}{d t}\left(\frac{\partial \mathcal{L}}{\partial \boldsymbol{\omega}}\right)+\tilde{v} \frac{\partial \mathcal{L}}{\partial \mathbf{v}}+\tilde{\omega} \frac{\partial \mathcal{L}}{\partial \boldsymbol{\omega}}-\left(E^{T}\right)^{-1} \frac{\partial \mathcal{L}}{\partial \boldsymbol{\theta}}\end{array} \\ & =\mathbf{M} \left\lvert\, \begin{array}{l}\text { Generalized } \\ \text { Forces }\end{array}\right. \\ \frac{\partial}{\partial t}\left(\frac{\partial \hat{\mathcal{L}}}{\partial \dot{\mathbf{u}}}\right)-\frac{\partial \hat{\mathcal{L}}}{\partial \mathbf{u}}+\frac{\partial \hat{\mathcal{R}}}{\partial \dot{\mathbf{u}}} & =\hat{\mathbf{U}}\end{aligned}$
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|  | Equations: onlinear ODEs | $\frac{d}{d t}\left(\frac{\partial \mathcal{L}}{\partial \mathbf{v}}\right)+\tilde{\omega} \frac{\partial \mathcal{L}}{\partial \mathbf{v}}-C \frac{\partial \mathcal{L}}{\partial \mathbf{R}}=$ | F |
| :---: | :---: | :---: | :---: |
| Moment Equations: 3 nonlinear ODEs | $\frac{d}{d t}\left(\frac{\partial \mathcal{L}}{\partial \boldsymbol{\omega}}\right)$ | $\tilde{v} \frac{\partial \mathcal{L}}{\partial \mathbf{v}}+\tilde{\omega} \frac{\partial \mathcal{L}}{\partial \boldsymbol{\omega}}-\left(E^{T}\right)^{-1} \frac{\partial \mathcal{L}}{\partial \boldsymbol{\theta}}=$ | M |
|  |  | $\frac{\partial}{\partial t}\left(\frac{\partial \hat{\mathcal{L}}}{\partial \dot{\mathbf{u}}}\right)-\frac{\partial \hat{\mathcal{L}}}{\partial \mathbf{u}}+\frac{\partial \hat{\mathcal{R}}}{\partial \dot{\mathbf{u}}}=$ | U |

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|  | Equations governing elastic displacement: 3 nonlinear PDEs | $\mathrm{g}: \frac{\partial}{\partial t}\left(\frac{\partial \hat{\mathcal{L}}}{\partial \dot{\mathbf{u}}}\right)-\frac{\partial \hat{\mathcal{L}}}{\partial \mathbf{u}}+\frac{\partial \hat{\mathcal{R}}}{\partial \dot{\mathbf{u}}}=$ | U |

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$(n+1)^{2}$ nodes .


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$\boldsymbol{\Delta}=\left[\begin{array}{lll}\mathbf{R}^{T} & \boldsymbol{\theta}^{T} & \mathbf{q}_{*}^{T}\end{array}\right]^{T} \quad$ Vector of nodal displacements
is the global displacement vector.

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Kinetic energy can be written in the quadratic form:

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is the global velocity vector.
Both $K$ and $M$ are $\left(n^{2}+4 n+6\right) \times\left(n^{2}+4 n+6\right)$.

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\mathbf{f}^{(i)}=2 \frac{S_{0}}{c}\left[\frac{R_{0}}{R^{(i)}}\right]^{2}\left[\hat{\mathbf{R}}^{(i)} \cdot \hat{\mathbf{n}}^{(i)}\right]^{2} \hat{\mathbf{n}}^{(i)}
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Unit vector along position vector from $O$ to a point in $i$ th finite element

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The solar radiation pressure over the $i$ th finite element is then

Solar radiation flux
Speed of light $\mathbf{f}^{(i)}=2 \frac{S_{0}}{c}\left[\frac{R_{0}}{R^{(i)}}\right]^{2}\left[\hat{\mathbf{R}}^{〔 i)} \cdot \hat{\mathbf{n}}^{(i)}\right]^{2} \hat{\mathbf{n}}_{\substack{(i)}}^{\begin{array}{l}\text { vector from } O \text { to a point in } i \text { th } \\ \text { finite element }\end{array}} \begin{aligned} & \text { Unit normal vector to } \\ & \text { ith element }\end{aligned}$

## EQUATIONS OF MOTION

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$\mathbf{f}^{(i)}=2 \frac{S_{0}}{c}\left[\frac{R_{0}}{R^{(i)}}\right]^{2}\left[\hat{\mathbf{R}}^{\zeta i)} \cdot \hat{\mathbf{n}}^{(i)}\right]^{2} \hat{\mathbf{n}}_{\leftarrow}^{(i)} \quad \begin{aligned} & \text { vector from } O \text { to a point in } i \text { th } \\ & \text { finite element }\end{aligned}$
Speed of light
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The solar radiation force over the control vanes are

## EQUATIONS OF MOTION

We assume that the sail film is perfect reflector.
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Solar radiation flux
$\underset{\text { of light }}{\mathbf{f}^{(i)}=2} \xrightarrow{S_{0}}\left[\frac{R_{0}}{R^{(i)}}\right]^{2}\left[\hat{\mathbf{R}}^{(i)} \cdot \hat{\mathbf{n}}^{(i)}\right]^{2} \hat{\mathbf{n}}^{(i)}$
Unit vector along position

Speed of light
The solar radiation force over the control vanes are

$$
\mathbf{f}_{c}^{(i)}=2 \frac{S_{0}}{c}\left[\frac{R_{0}}{R_{c}^{(i)}}\right]^{2} l^{2}\left[\hat{\mathbf{R}}_{c}^{(i)} \cdot \hat{\mathbf{n}}_{c}^{(i)}\right]^{2} \hat{\mathbf{n}}_{c}^{(i)}, \quad i=1,2,3,4
$$

## EQUATIONS OF MOTION

The total virtual work is

$$
\delta W=\delta W_{m}+\delta W_{b}+\delta W_{c}+\delta W_{p}=\mathbf{F}^{T} \delta \mathbf{R}^{*}+\mathbf{M}^{T} \delta \boldsymbol{\theta}^{*}+\mathbf{Q}^{T} \delta \mathbf{q}
$$

## EQUATIONS OF MOTION

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$\dot{\mathbf{x}}(t)=\mathbf{f}[\mathbf{x}(t), \mathbf{u}(t)]$
where
$\mathbf{x}(t)=\left[\begin{array}{c}\boldsymbol{\Delta}(t) \\ \mathbf{V}(t)\end{array}\right], \quad \mathbf{u}(t)=\left[\begin{array}{c}\beta_{1}(t) \\ \vdots \\ \beta_{4}(t) \\ \delta_{1}(t) \\ \vdots \\ \delta_{4}(t)\end{array}\right]$

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## EQUATIONS OF MOTION

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- gravitational forces


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Due to the nature of FEM, a large $n$ must be used for sufficiently accurate representation of the system.

## EQUATIONS OF MOTION

The system is also underactuated because the number of control inputs is many times smaller than the number of degrees of freedom.

## PERTURBATION SOLUTION

The problem with the nonlinearity of the system can be obviated by a perturbation solution in which

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\mathbf{x}(t)=\overline{\mathbf{x}}(t)+\hat{\mathbf{x}}(t)
$$

$$
\mathbf{u}(t)=\overline{\mathbf{u}}(t)+\hat{\mathbf{u}}(t)
$$

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for the desired nominal dynamics

## PERTURBATION SOLUTION

The perturbation solution separates the equation into the zero-order equation

$$
\dot{\overline{\mathbf{x}}}(t)=\mathbf{f}[\overline{\mathbf{x}}(t), \overline{\mathbf{u}}(t)]
$$

for the desired nominal dynamics and the first-order equation
$\dot{\hat{\mathbf{x}}}(t)=A[\overline{\mathbf{x}}(t), \overline{\mathbf{u}}(t)] \hat{\mathbf{x}}(t)+B[\overline{\mathbf{x}}(t), \overline{\mathbf{u}}(t)] \hat{\mathbf{u}}(t)$
for the small perturbation about the nominal dynamics.

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$\dot{\overline{\mathbf{x}}}(t)=\mathbf{f}[\overline{\mathbf{x}}(t), \overline{\mathbf{u}}(t)]$
for the desired nominal dynamics and the first-order equation
$\dot{\hat{\mathbf{x}}}(t)=A^{A}[\overline{\mathbf{x}}(t), \overline{\mathbf{u}}(t)] \hat{\mathbf{x}}(t)+B[\overline{\mathbf{x}}(t), \overline{\mathbf{u}}(t)] \hat{\mathbf{u}}(t)$
for the small perturbation about the nominal dynamics.

Coefficient matrices, both functions of $\overline{\mathbf{x}}(t)$ and $\overline{\mathbf{u}}(t)$

## PERTURBATION SOLUTION

The zero-order equation is highly nonlinear and high dimensional.

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Its inverse dynamics is used to compute an open-loop control $\overline{\mathbf{u}}(t)$ to achieve a desired zero-order state $\overline{\mathbf{x}}(t)$.

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The first-order equation is linear, but as high dimensional as the original equation.

## PERTURBATION SOLUTION

The zero-order equation is highly nonlinear and high dimensional.
Its inverse dynamics is used to compute an open-loop control $\overline{\mathbf{u}}(t)$ to achieve a desired zero-order state $\overline{\mathbf{x}}(t)$.

The first-order equation is linear, but as high dimensional as the original equation.

We note that the zero-order state $\overline{\mathbf{x}}(t)$ and control input $\overline{\mathbf{u}}(t)$ enter into the first-order equation as inputs.

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As a result, the first-order equation is time-invariant if $\overline{\mathbf{x}}$ and $\overline{\mathbf{u}}$ are constant.

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The first-order equation is used to assess stability about the maneuver.

## PERTURBATION SOLUTION

As a result, the first-order equation is time-invariant if $\overline{\mathbf{x}}$ and $\overline{\mathbf{u}}$ are constant.

The first-order equation is time-varying if $\overline{\mathbf{x}}$ and $\overline{\mathbf{u}}$ are time-dependent.
The first-order equation is used to assess stability about the maneuver.
It can also be used to design feedback control to attenuate perturbations.

## MODEL REDUCTION

The problems with high-dimensionality and underactuation are obviated by a model-reduction approach in which

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$$

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$$

- Vector of nodal displacements


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$$
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$$

where $U_{e}$ is the matrix of $k$ vibration modes and
$\boldsymbol{\xi}=\left[\begin{array}{c}\xi_{1} \\ \xi_{2} \\ \vdots \\ \xi_{k}\end{array}\right], \quad \boldsymbol{\eta}=\dot{\boldsymbol{\xi}}=\left[\begin{array}{c}\eta_{1} \\ \eta_{2} \\ \vdots \\ \eta_{k}\end{array}\right]$

## NUMERICAL APPLICATION

## Numerical Data

| Length, $L=100 \mathrm{~m}$ | Boom length, $L_{b}=100 \sqrt{2} \mathrm{~m}$ |
| :---: | :---: |
| Membrane thickness, $t_{m}=2.5 \mu \mathrm{~m}$ | Membrane density, $\rho_{m}=1660 \mathrm{~kg} / \mathrm{m}^{3}$ |
| Boom wall thickness, $t_{b}=0.1 \mathrm{~mm}$ | Boom Radius, $r_{b}=3.5 \mathrm{~cm}$ |
| Boom density, $\rho_{b}=1660 \mathrm{~kg} / \mathrm{m}^{3}$ | Boom Young's Modulus, $E_{b}=68.95 \mathrm{GPa}$ |
| Payload mass, $M_{p}=20 \mathrm{~kg}$ | Payload position, $\mathbf{r}_{p}=\left[\begin{array}{lll}0 & 0 & -0.1\end{array}\right]^{T} \mathrm{~m}$ |
| Control panel length, $l=5 \mathrm{~m}$ | Control panel mass $m_{c}=0.3124 \mathrm{~kg}$ |
| Damping factor $\zeta=0.005$ | Tension, $T=0.0172 \mathrm{~N} / \mathrm{m}$ |
| Solar Radiation flux, $S_{0}=1368 \mathrm{~N} / \mathrm{m}^{2}$ | Sun to Earth Distance, $R_{0}=1.496 \times 10^{11} \mathrm{~m}$ |

## NUMERICAL APPLICATION

First 4 nonzero eigenfrequencies vs. $n$





## NUMERICAL APPLICATION

Shapes of first 8 vibration modes for $n=14$

7 th Mode at $\omega_{7}=0.092156 \mathrm{rad} / \mathrm{s}$


8th Mode at $\omega_{8}=0.104405 \mathrm{rad} / \mathrm{s}$

9th Mode at $\omega_{9}=0.128955 \mathrm{rad} / \mathrm{s}$



10th Mode at $\omega_{10}=0.144522 \mathrm{rad} / \mathrm{s}$


## NUMERICAL APPLICATION

Shapes of first 8 vibration modes for $n=14$ (continued)

11th Mode at $\omega_{11}=0.144522 \mathrm{rad} / \mathrm{s}$


12th Mode at $\omega_{12}=0.164341 \mathrm{rad} / \mathrm{s}$


14th Mode at $\omega_{14}=0.220341 \mathrm{rad} / \mathrm{s}$


## NUMERICAL APPLICATION

Solarcraft in a circular orbit around the Sun at $R=R_{0}=1 \mathrm{AU}$


## NUMERICAL APPLICATION

Solarcraft in a circular orbit around the Sun at $R=R_{0}=1 \mathrm{AU}$


## NUMERICAL APPLICATION

Solarcraft in a circular orbit around the Sun at $R=R_{0}=1 \mathrm{AU}$


## NUMERICAL APPLICATION

Convergence of the elastic displacements at the points $a$ and $b$


## NUMERICAL APPLICATION



## NUMERICAL APPLICATION

Elastic shape for $n=14$ during the circular orbit


## NUMERICAL APPLICATION

## Equilibrium Values and Eigenvalues of $A$ for the circular orbit at 1 AU

|  | $n=6$ | $n=10$ | $n=14$ |
| :--- | :--- | :--- | :--- |
| $\Omega$ | $1.767 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $1.767 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $1.767 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ |
| $w_{a}$ | -1.4668 m | -1.4479 m | -1.4444 m |
| $w_{b}$ | -0.9380 m | -0.9233 m | -0.9209 m |
| $\#$ of $\lambda$ 's | 132 | 292 | 516 |
| $\lambda_{1-4}$ | 0 | 0 | 0 |
| $\lambda_{5,6}$ | $\pm 1.438 \times 10^{-7}$ | $\pm 1.416 \times 10^{-7}$ | $\pm 1.355 \times 10^{-7}$ |
| $\lambda_{7,8}$ | $\pm i 1.767 \times 10^{-7}$ | $\pm i 1.766 \times 10^{-7}$ | $\pm i 1.769 \times 10^{-7}$ |
| $\lambda_{9,10}$ | $\pm i 2.421 \times 10^{-7}$ | $\pm i 2.388 \times 10^{-7}$ | $\pm i 2.685 \times 10^{-7}$ |
| $\lambda_{11,12}$ | $\pm 0.001148$ | $\pm 0.001112$ | $\pm 0.001101$ |
| $\lambda_{13,14}$ | $-0.000332 \pm i 0.078881$ | $-0.000332 \pm i 0.078351$ | $-0.000332 \pm i 0.078170$ |
| $\lambda_{15,16}$ | $-0.000469 \pm i 0.093866$ | $-0.000463 \pm i 0.092554$ | $-0.000461 \pm i 0.092144$ |
| $\lambda_{17,18}$ | $-0.000943 \pm i 0.133046$ | $-0.000910 \pm i 0.129813$ | $-0.000902 \pm i 0.128940$ |
| $\lambda_{19,20}$ | $-0.001199 \pm i 0.150027$ | $-0.001148 \pm i 0.145763$ | $-0.001133 \pm i 0.144508$ |
| $\lambda_{21,22}$ | $-0.001212 \pm i 0.150838$ | $-0.001159 \pm i 0.146498$ | $-0.001145 \pm i 0.145234$ |
| $\lambda_{23,24}$ | $-0.001450 \pm i 0.164970$ | $-0.001385 \pm i 0.160137$ | $-0.001365 \pm i 0.158598$ |

## NUMERICAL APPLICATION

## Equilibrium Values and Eigenvalues of $A$ for the circular orbit at 1 AU

| $\lambda_{25,26}$ | $-0.002715 \pm i 0.225751$ | $-0.002405 \pm i 0.210986$ | $-0.002320 \pm i 0.206787$ |
| :--- | :--- | :--- | :--- |
| $\lambda_{27,28}$ | $-0.003044 \pm i 0.239051$ | $-0.002731 \pm i 0.224812$ | $-0.002634 \pm i 0.220319$ |
| $\lambda_{29,30}$ | $-0.003045 \pm i 0.239097$ | $-0.002733 \pm i 0.224893$ | $-0.002636 \pm i 0.220414$ |
| $\lambda_{31,32}$ | $-0.003383 \pm i 0.251986$ | $-0.003087 \pm i 0.239036$ | $-0.002980 \pm i 0.234326$ |
| $\lambda_{33,34}$ | $-0.004655 \pm i 0.295580$ | $-0.003896 \pm i 0.268537$ | $-0.003689 \pm i 0.260732$ |
| $\lambda_{35,36}$ | $-0.004670 \pm i 0.296069$ | $-0.003957 \pm i 0.270620$ | $-0.003765 \pm i 0.263389$ |
| $\lambda_{37,38}$ | $-0.004759 \pm i 0.298872$ | $-0.004083 \pm i 0.274897$ | $-0.003882 \pm i 0.267456$ |
| $\lambda_{39,40}$ | $-0.004801 \pm i 0.300184$ | $-0.004094 \pm i 0.275274$ | $-0.003890 \pm i 0.267715$ |
| $\lambda_{41,42}$ | $-0.007565 \pm i 0.376808$ | $-0.005233 \pm i 0.311208$ | $-0.004876 \pm i 0.299754$ |
| $\lambda_{43,44}$ | $-0.007847 \pm i 0.383764$ | $-0.005298 \pm i 0.313130$ | $-0.004891 \pm i 0.300194$ |
| $\lambda_{45,46}$ | $-0.008280 \pm i 0.394194$ | $-0.005363 \pm i 0.315035$ | $-0.004985 \pm i 0.303071$ |
| $\lambda_{47,48}$ | $-0.009178 \pm i 0.415023$ | $-0.005558 \pm i 0.320716$ | $-0.005214 \pm i 0.309937$ |
| $\lambda_{49,50}$ | $-0.009414 \pm i 0.420309$ | $-0.006960 \pm i 0.358878$ | $-0.006383 \pm i 0.342945$ |
| $\lambda_{51,52}$ | $-0.009545 \pm i 0.423216$ | $-0.007144 \pm i 0.363611$ | $-0.006547 \pm i 0.347303$ |
| $\lambda_{53,54}$ | $-0.009557 \pm i 0.423487$ | $-0.007166 \pm i 0.364171$ | $-0.006559 \pm i 0.347639$ |
| $\lambda_{55,56}$ | $-0.009557 \pm i 0.423487$ | $-0.007224 \pm i 0.365625$ | $-0.006633 \pm i 0.349576$ |
| $\lambda_{57,58}$ | $-0.009557 \pm i 0.423487$ | $-0.008991 \pm i 0.407885$ | $-0.008106 \pm i 0.386430$ |

## NUMERICAL APPLICATION

## Equilibrium Values and Eigenvalues of $A$ for $n=14$

|  | $0.25 R_{0}$ | $0.5 R_{0}$ | $0.75 R_{0}$ | $1.25 R_{0}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\Omega$ | $9.615 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $4.999 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $3.899 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $2.5284 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ |
| $w_{a}$ | -4.9852 m | -5.7761 m | -6.5373 m | -5.1013 m |
| $w_{b}$ | -2.8067 m | -3.6829 m | -4.0839 m | -2.4573 m |
| $\lambda_{1-4}$ | 0 | 0 | 0 | 0 |
| $\lambda_{5,6}$ | $\pm 7.286 \times 10^{-8}$ | $\pm 2.070 \times 10^{-7}$ | $\pm 2.960 \times 10^{-7}$ | $\pm 4.370 \times 10^{-7}$ |
| $\lambda_{7,8}$ | $\pm i 9.614 \times 10^{-7}$ | $\pm i 4.997 \times 10^{-7}$ | $\pm i 3.899 \times 10^{-7}$ | $\pm i 2.529 \times 10^{-7}$ |
| $\lambda_{9,10}$ | $\pm i 1.243 \times 10^{-6}$ | $\pm i 8.798 \times 10^{-7}$ | $\pm i 7.670 \times 10^{-7}$ | $\pm i 6.161 \times 10^{-7}$ |
| $\lambda_{11,12}$ | $\pm 0.001252$ | $\pm 0.004405$ | $\pm 0.007057$ | $\pm 0.010065$ |
| $\lambda_{13,14}$ | $-0.000332 \pm i 0.078184$ | $-0.000332 \pm i 0.078184$ | $-0.000332 \pm i 0.078184$ | $-0.000332 \pm i 0.078184$ |
| $\lambda_{15,16}$ | $-0.000461 \pm i 0.092155$ | $-0.000461 \pm i 0.092155$ | $-0.000461 \pm i 0.092155$ | $-0.000461 \pm i 0.092155$ |
| $\lambda_{17,18}$ | $-0.000902 \pm i 0.128951$ | $-0.000902 \pm i 0.128951$ | $-0.000902 \pm i 0.128951$ | $-0.000902 \pm i 0.128951$ |
| $\lambda_{19,20}$ | $-0.001133 \pm i 0.144518$ | $-0.001133 \pm i 0.144514$ | $-0.001133 \pm i 0.144510$ | $-0.001133 \pm i 0.144516$ |
| $\lambda_{21,22}$ | $-0.001145 \pm i 0.145244$ | $-0.001145 \pm i 0.145244$ | $-0.001145 \pm i 0.145244$ | $-0.001145 \pm i 0.145244$ |
| $\lambda_{23,24}$ | $-0.001365 \pm i 0.158607$ | $-0.001365 \pm i 0.158607$ | $-0.001365 \pm i 0.158607$ | $-0.001365 \pm i 0.158607$ |
| $\lambda_{25,26}$ | $-0.002320 \pm i 0.206795$ | $-0.002320 \pm i 0.206795$ | $-0.002320 \pm i 0.206795$ | $-0.002320 \pm i 0.206795$ |
| $\lambda_{27,28}$ | $-0.002634 \pm i 0.220326$ | $-0.002634 \pm i 0.220326$ | $-0.002634 \pm i 0.220327$ | $-0.002634 \pm i 0.220333$ |
| $\lambda_{29,30}$ | $-0.002636 \pm i 0.220421$ | $-0.002636 \pm i 0.220421$ | $-0.002636 \pm i 0.220421$ | $-0.002636 \pm i 0.220421$ |
| $\lambda_{31,32}$ | $-0.002980 \pm i 0.234332$ | $-0.002980 \pm i 0.234332$ | $-0.002980 \pm i 0.234332$ | $-0.002980 \pm i 0.234332$ |
| $\lambda_{33,34}$ | $-0.003689 \pm i 0.260738$ | $-0.003689 \pm i 0.260738$ | $-0.003689 \pm i 0.260738$ | $-0.003689 \pm i 0.260738$ |
| $\lambda_{35,36}$ | $-0.003765 \pm i 0.263395$ | $-0.003765 \pm i 0.263395$ | $-0.003765 \pm i 0.263395$ | $-0.003765 \pm i 0.263395$ |
| $\vdots$ | $\vdots$ | $\vdots$ | $\vdots$ | $\vdots$ |
|  |  |  |  |  |

## NUMERICAL APPLICATION

Elastic shapes for $n=14$


## NUMERICAL APPLICATION

## Equilibrium Values and Eigenvalues of $A_{r}$ for $n=14$

|  | $k=8$ | $k=8$ (Closed-Loop) | $k=10$ | $k=12$ |
| :--- | :--- | :--- | :--- | :--- |
| $\Omega$ | $1.767 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $1.767 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $1.767 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ | $1.767 \times 10^{-7} \mathrm{rad} / \mathrm{s}$ |
| $w_{a}$ | -1.4393 m | -1.4393 m | -1.4393 m | -1.4393 m |
| $w_{b}$ | -0.9493 m | -0.9493 m | -0.9493 m | -0.9493 m |
| $\lambda_{1-4}$ | 0 | 0 | 0 | 0 |
| $\lambda_{5,6}$ | $\pm 8.145 \times 10^{-8}$ | $3.31 \times 10^{-6} \pm i 0.000029$ | $\pm 1.451 \times 10^{-7}$ | $\pm 1.451 \times 10^{-7}$ |
| $\lambda_{7,8}$ | $\pm i 1.767 \times 10^{-7}$ | $\pm i 1.767 \times 10^{-7}$ | $\pm i 1.767 \times 10^{-7}$ | $\pm i 1.767 \times 10^{-7}$ |
| $\lambda_{9,10}$ | $\pm i 3.096 \times 10^{-7}$ | $-0.076316 \pm i 0.076471$ | $\pm i 2.432 \times 10^{-7}$ | $\pm i 2.432 \times 10^{-7}$ |
| $\lambda_{11,12}$ | $\pm 0.001171$ | $-0.001431 \pm i 0.001351$ | $\pm 0.001171$ | $\pm 0.001171$ |
| $\lambda_{13,14}$ | $-0.000461 \pm i 0.092155$ | $-0.000461 \pm i 0.092155$ | $-0.000461 \pm i 0.092144$ | $-0.000461 \pm i 0.092144$ |
| $\lambda_{15,16}$ | $-0.000902 \pm i 0.128951$ | $-0.000902 \pm i 0.128951$ | $-0.000902 \pm i 0.128937$ | $-0.000902 \pm i 0.128937$ |
| $\lambda_{17,18}$ | $-0.001106 \pm i 0.142740$ | $-0.001106 \pm i 0.142740$ | $-0.001105 \pm i 0.142725$ | $-0.001105 \pm i 0.142716$ |
| $\lambda_{19,20}$ | $-0.001133 \pm i 0.144518$ | $-0.001251 \pm i 0.144029$ | $-0.001133 \pm i 0.144506$ | $-0.001133 \pm i 0.144506$ |
| $\lambda_{21,22}$ | $-0.001372 \pm i 0.159001$ | $-0.001372 \pm i 0.159001$ | $-0.001372 \pm i 0.158991$ | $-0.001372 \pm i 0.158991$ |
| $\lambda_{23,24}$ | $-0.002320 \pm i 0.206795$ | $-0.002320 \pm i 0.206795$ | $-0.002320 \pm i 0.206785$ | $-0.002320 \pm i 0.206785$ |
| $\lambda_{25,26}$ | $-0.002626 \pm i 0.219981$ | $-0.002627 \pm i 0.219961$ | $-0.002618 \pm i 0.219630$ | $-0.002617 \pm i 0.219613$ |
| $\lambda_{27,28}$ | $-0.000721 \pm i 1281.120$ | $-0.000721 \pm i 1281.120$ | $-0.002634 \pm i 0.220317$ | $-0.002634 \pm i 0.220317$ |
| $\lambda_{29,30}$ |  |  | $-0.002980 \pm i 0.234325$ | $-0.002980 \pm i 0.234325$ |
| $\lambda_{31,32}$ |  |  | $-0.000729 \pm i 1283.280$ | $-0.003765 \pm i 0.263390$ |
| $\lambda_{33,34}$ |  |  |  | $-0.003836 \pm i 0.265872$ |
| $\lambda_{35,36}$ |  |  | $-0.000776 \pm i 1291.120$ |  |

## NUMERICAL APPLICATION

Equilibrium values for $n=14$ at a circular orbit at the L1 libration $\operatorname{point}\left(\beta_{1}=\beta_{2}=\beta_{3}=\beta_{4}=0, \delta_{1}=\delta_{2}=\delta_{3}=\delta_{4}=\delta\right.$ )

|  | $\delta=0 \mathrm{rad}$ | $\delta=\pi / 6 \mathrm{rad}$ | $\delta=\pi / 4 \mathrm{rad}$ | $\delta=\pi / 3 \mathrm{rad}$ |
| :--- | :--- | :--- | :--- | :--- |
| $R$ | $1.38144 \times 10^{11} \mathrm{~m}$ | $1.38187 \times 10^{11} \mathrm{~m}$ | $1.38224 \times 10^{11} \mathrm{~m}$ | $1.38252 \times 10^{11} \mathrm{~m}$ |
| $R_{0}-R$ | $1.14555 \times 10^{10} \mathrm{~m}$ | $1.14126 \times 10^{10} \mathrm{~m}$ | $1.13763 \times 10^{10} \mathrm{~m}$ | $1.13484 \times 10^{10} \mathrm{~m}$ |
| $w_{a}$ | -1.6924 m | -1.6977 m | -1.7022 m | -1.7056 m |
| $w_{b}$ | -1.0791 m | -1.0806 m | -1.0819 m | -1.0829 m |

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For a conventional spacecraft $R=1.48118 \times 10^{11} \mathrm{~m}$ and $R-R_{0}=1.48242 \times 10^{9} \mathrm{~m}$.

## CONCLUSIONS

This presents a comprehensive mathematical model for the dynamics of a solarcraft, accounting for both rigid body and elastic motions and their interactions.

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The zero-order equation is used to design an open-loop control for a desired maneuver.

## CONCLUSIONS

The first-order equation is used to assess the stability about the desired maneuver and to design feedback control to alleviate perturbations.

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High dimensionality and underactuation are dealt with a model reduction approach.

In numerical application, circular orbits around the Sun are considered. The zero-order equation is used to determine the turn rate and elastic deformation.

## CONCLUSIONS

Circular orbit at the sub-L1 libration point is also considered.
The first-order equation is used to assess the stability about the orbits.

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The first-order equation is used to assess the stability about the orbits.
Future research will focus on introduction of more effective control inputs, and control design.

